

1. Calculate the first variation of the following functionals:

$$a. A(u) = \int_0^1 x u(x)^2 + \left( \frac{\partial u(x)}{\partial x} \right)^2 dx$$

Penyelesaian:

$$\begin{aligned} \frac{\partial}{\partial \varepsilon} A(\psi + \varepsilon \eta) \Big|_{\varepsilon=0} &= \frac{\partial}{\partial \varepsilon} \int_0^1 x (\psi + \varepsilon \eta)^2 + \left( \frac{\partial (\psi + \varepsilon \eta)}{\partial x} \right)^2 dx \Big|_{\varepsilon=0} \\ &= \frac{\partial}{\partial \varepsilon} \int_0^1 x (\psi^2 + 2u\varepsilon\eta + \varepsilon^2\eta^2) + (\partial_x \psi + \varepsilon \partial_x \eta)^2 dx \Big|_{\varepsilon=0} \\ &= \int_0^1 2u\eta x dx + \int_0^1 2\partial_x u \partial_x \eta dx \\ &= \int_0^1 2u\eta x dx + 2\partial_x u \eta - \int_0^1 2\partial_{xx} u \eta dx \\ &= 2\partial_x u \eta \Big|_0^1 + \int_0^1 (\psi u x - 2\partial_{xx} u) \eta dx \quad \Leftrightarrow \text{Bcs} + \langle \delta A, \eta \rangle \end{aligned}$$

Maka turunan variasi dari  $A(\psi(x))$  adalah  $\delta A(\psi(x)) = 2xu(x) - 2\partial_{xx}u(x)$ .

$$b. A(u) = \int \sin x u(x)^2 + x^3 \left( \frac{\partial u(x)}{\partial x} \right)^2 dx$$

Penyelesaian:

$$\begin{aligned} \frac{\partial}{\partial \varepsilon} A(\psi + \varepsilon \eta) \Big|_{\varepsilon=0} &= \frac{\partial}{\partial \varepsilon} \int_0^1 \sin x (\psi + \varepsilon \eta)^2 + x^3 \left( \frac{\partial (\psi + \varepsilon \eta)}{\partial x} \right)^2 dx \Big|_{\varepsilon=0} \\ &= \frac{\partial}{\partial \varepsilon} \int_0^1 \sin x (\psi^2 + 2u\varepsilon\eta + \varepsilon^2\eta^2) + x^3 (\partial_x \psi + \varepsilon \partial_x \eta)^2 dx \Big|_{\varepsilon=0} \\ &= \int_0^1 2u\eta \sin x dx + \int_0^1 2x^3 \partial_x u \partial_x \eta dx \\ &= \int_0^1 2u\eta \sin x dx + 2x^3 \partial_x u \eta \Big|_0^1 - \int_0^1 (x^2 \partial_x u + 2x^3 \partial_{xx} u) \eta dx \\ &= 2x^3 \partial_x u \eta \Big|_0^1 + \int_0^1 (\psi u \sin x - 6x^2 \partial_x u - 2x^3 \partial_{xx} u) \eta dx \quad \Leftrightarrow \text{Bcs} + \langle \delta A, \eta \rangle \end{aligned}$$

Maka turunan variasi dari  $A \underline{\Phi}(x)$  adalah

$$\delta A \underline{\Phi}(x) = 2u(x)\sin x - 6x^2\partial_x u(x) - 2x^3\partial_{xx} u(x).$$

c.  $A(u) = \int u(x)^2 + \left( \frac{\partial u(x)}{\partial x} \right)^2 dx$

Penyelesaian:

$$\begin{aligned} \frac{\partial}{\partial \varepsilon} A \underline{\Phi} + \varepsilon \eta \Big|_{\varepsilon=0} &= \frac{\partial}{\partial \varepsilon} \int \underline{\Phi} + \varepsilon \eta \Big|_{\varepsilon=0} + \left( \frac{\partial \underline{\Phi} + \varepsilon \eta}{\partial x} \right)^2 dx \Big|_{\varepsilon=0} \\ &= \int \frac{\partial}{\partial \varepsilon} \underline{\Phi} + \varepsilon \eta \Big|_{\varepsilon=0} + \frac{\partial}{\partial \varepsilon} \Phi_x u + \varepsilon \partial_x \eta \Big|_{\varepsilon=0} dx \\ &= \int 2u\eta dx + \int 2\partial_x u \partial_x \eta dx \\ &= \int 2u\eta dx + 2\Phi_x u \eta \Big|_{x=D} - \int 2\partial_{xx} u \eta dx \\ &= 2\partial_x u \eta \Big|_{x=D} + \int \Phi u - 2\partial_{xx} u \eta dx \\ &\Leftrightarrow \text{Bcs} + \langle \delta A, \eta \rangle \end{aligned}$$

Turunan variasi dari  $A \underline{\Phi}(x)$  adalah  $\delta A \underline{\Phi}(x) = 2u(x) - 2\partial_{xx} u(x)$ .

d.  $A(u) = \int_0^1 \sin u(x) + \left( \frac{\partial^2 u(x)}{\partial x^2} \right)^2 dx$

Penyelesaian:

$$\begin{aligned} \frac{\partial}{\partial \varepsilon} A \underline{\Phi} + \varepsilon \eta \Big|_{\varepsilon=0} &= \frac{\partial}{\partial \varepsilon} \int_0^1 \sin \underline{\Phi} + \varepsilon \eta \Big|_{\varepsilon=0} + \left( \frac{\partial^2 \underline{\Phi} + \varepsilon \eta}{\partial x^2} \right)^2 dx \Big|_{\varepsilon=0} \\ &= \frac{\partial}{\partial \varepsilon} \int_0^1 \sin \underline{\Phi} + \varepsilon \eta \Big|_{\varepsilon=0} + \Phi_x (\partial_x u + \varepsilon \partial_x \eta) \Big|_{\varepsilon=0} dx \\ &= \frac{\partial}{\partial \varepsilon} \int_0^1 \sin \underline{\Phi} + \varepsilon \eta \Big|_{\varepsilon=0} + \Phi_{xx} u + \varepsilon \partial_{xx} \eta \Big|_{\varepsilon=0} dx \\ &= \int_0^1 \cos(u)\eta dx + \int_0^1 2\Phi_{xx} u \partial_{xx} \eta dx \end{aligned}$$

$$\begin{aligned}
&= \int_0^1 \cos(u) \eta \, dx + 2 \oint_{xx} u \partial_x \eta \Big|_0^1 - \int_0^1 2 \oint_{xxx} u \partial_x \eta \, dx \\
&= 2 \oint_{xx} u \partial_x \eta \Big|_0^1 - 2 \oint_{xxx} u \eta \Big|_0^1 + \int_0^1 \cos(u) \eta \, dx + \int_0^1 2 \oint_x^{(iv)} u \eta \, dx \\
&= 2 \oint_{xx} u \partial_x \eta \Big|_0^1 - 2 \oint_{xxx} u \eta \Big|_0^1 + \int_0^1 (\cos(u) + 2 \oint_x^{(iv)} u) \eta \, dx
\end{aligned}$$

Maka turunan variasi dari  $A(\psi(x))$  adalah  $\delta A(\psi(x)) = \cos(u(x)) + 2 \partial_x^{(iv)} u(x)$ .

e.  $A(u) = \int_0^1 u(x)^4 + \left( \frac{\partial u(x)}{\partial x} \right)^7 \, dx$

Penyelesaian:

$$\begin{aligned}
\frac{\partial}{\partial \varepsilon} A(\psi + \varepsilon \eta) \Big|_{\varepsilon=0} &= \frac{\partial}{\partial \varepsilon} \int_0^1 (\psi + \varepsilon \eta)^4 + \left( \frac{\partial (\psi + \varepsilon \eta)}{\partial x} \right)^7 \, dx \Big|_{\varepsilon=0} \\
&= \int_0^1 \frac{\partial}{\partial \varepsilon} (\psi + \varepsilon \eta)^4 \Big|_{\varepsilon=0} + \frac{\partial}{\partial \varepsilon} \oint_x u + \varepsilon \partial_x \eta \Big|_{\varepsilon=0} \, dx \\
&= \int_0^1 4u^3 \eta \, dx + \int_0^1 7 \oint_x u \partial_x \eta \, dx \\
&= \int_0^1 4u^3 \eta \, dx + 7 \oint_x u \eta \Big|_0^1 - \int_0^1 42 \oint_x u \partial_{xx} u \eta \, dx \\
&= 7 \oint_x u \eta \Big|_0^1 + \int_0^1 (4u^3 - 42 \oint_x u \partial_{xx} u) \eta \, dx \\
&\Leftrightarrow \text{Bcs} + \langle \delta A, \eta \rangle
\end{aligned}$$

Turunan variasi dari  $A(\psi(x))$  adalah  $\delta A(\psi(x)) = 4\psi(x)^3 - 42 \oint_x u(x) \partial_{xx} u(x)$ .

f.  $A(u) = \int_0^1 n(x) \left( 1 + \frac{\partial u(x)^2}{\partial x} \right)^{\frac{1}{2}} \, dx$

Penyelesaian:

Persamaan untuk turunan variasi adalah

$$\delta A(u) = \frac{\partial L}{\partial u} - \frac{\partial}{\partial x} \frac{\partial L}{\partial \dot{u}}$$

dimana  $L$  adalah *integrand* dari *functional* dan  $\dot{u} = \frac{\partial u(x)}{\partial x}$ . Maka turunan variasi

untuk soal no 1.f adalah

$$\begin{aligned}\delta A \underbrace{L}_{\text{dengan}} &= \frac{\partial}{\partial x} \left[ \frac{\underbrace{L(x) \partial_x u(x)}_{\text{dengan}}}{\sqrt{1 + \underbrace{\partial_x u(x)}^2}} \right] \\ &= \frac{\underbrace{\partial_x n(x) \partial_x u(x)}_{\text{dengan}}}{\sqrt{1 + \underbrace{\partial_x u(x)}^2}} + \frac{n(x) \underbrace{\partial_{xx} u(x)}_{\text{dengan}}}{\sqrt{1 + \underbrace{\partial_x u(x)}^2}} - \frac{n(x) \underbrace{\partial_x u(x)}^2 \underbrace{\partial_{xx} u(x)}_{\text{dengan}}}{(1 + \underbrace{\partial_x u(x)}^2)^{3/2}}\end{aligned}$$

g.  $A(q) = \int \frac{1}{2} \dot{q}(t)^2 - \frac{1}{2} q(t)^2 + q(t)^3 dt$

Penyelesaian:

$$\begin{aligned}\frac{\partial}{\partial \varepsilon} A \underbrace{q + \varepsilon \eta}_{\text{dengan}} \Big|_{\varepsilon=0} &= \frac{\partial}{\partial \varepsilon} \int \frac{1}{2} \left( \frac{\partial \underbrace{q + \varepsilon \eta}_{\text{dengan}}}{\partial t} \right)^2 - \frac{1}{2} \underbrace{q + \varepsilon \eta}_{\text{dengan}}^2 + \underbrace{q + \varepsilon \eta}_{\text{dengan}}^3 dt \Big|_{\varepsilon=0} \\ &= \int \frac{\partial}{\partial \varepsilon} \left( \frac{1}{2} \underbrace{q(t) + \varepsilon \dot{q}(t)}_{\text{dengan}}^2 \right) \Big|_{\varepsilon=0} - \frac{\partial}{\partial \varepsilon} \left( \frac{1}{2} \underbrace{q(t) + \varepsilon \eta(t)}_{\text{dengan}}^2 \right)^2 \Big|_{\varepsilon=0} + \frac{\partial}{\partial \varepsilon} \underbrace{q(t) + \varepsilon \eta(t)}_{\text{dengan}}^3 \Big|_{\varepsilon=0} dt \\ &= \int \dot{q}(t) \dot{\eta}(t) dt - \int q(t) \eta(t) dt + \int 3 \underbrace{q(t)}^2 \eta(t) dt \\ &= \dot{q}(t) \eta(t) \Big|_{x=D} - \int \ddot{q}(t) \eta(t) dt - \int q(t) \eta(t) dt + \int 3 \underbrace{q(t)}^2 \eta(t) dt \\ &= \dot{q}(t) \eta(t) \Big|_{x=D} + \int (\underbrace{q(t)}^2 - \dot{q}(t) - q(t)) \dot{\eta}(t) dt \\ &\Leftrightarrow \text{Bcs} + \langle \delta A, \eta \rangle\end{aligned}$$

Turunan variasi dari  $A(q)$  adalah  $\delta A(q) = 3 \underbrace{q(t)}^2 - \dot{q}(t) - q(t)$

h.  $A(u) = \int \frac{1}{2} \partial_x u(x)^2 + x^3 \sin u(x) + u(x)^5 dx$

Penyelesaian:

$$\begin{aligned}
\frac{\partial}{\partial \varepsilon} A(\psi + \varepsilon\eta) \Big|_{\varepsilon=0} &= \frac{\partial}{\partial \varepsilon} \int \frac{1}{2} \left( \frac{\partial \psi + \varepsilon\eta}{\partial x} \right)^2 + x^3 \sin \psi + \varepsilon\eta \, dx \Big|_{\varepsilon=0} \\
&= \int \frac{1}{2} \frac{\partial}{\partial \varepsilon} (\psi_x u + \varepsilon \partial_x \eta) \Big|_{\varepsilon=0} + \frac{\partial}{\partial \varepsilon} x^3 \sin \psi + \varepsilon\eta \Big|_{\varepsilon=0} \, dx \\
&= \int \psi_x u \eta \, dx + \int x^3 \cos(u) \eta \, dx + \int 5u^4 \eta \, dx \\
&= \psi_x u \eta \Big|_{x=D} - \int \psi_{xx} u \eta \, dx + \int (\psi^3 \cos(u) + 5u^4) \eta \, dx \\
&= \psi_x u \eta \Big|_{x=D} + \int (\psi^3 \cos(u) + 5u^4 - \partial_{xx} u) \eta \, dx \\
&\Leftrightarrow \text{Bcs} + \langle \delta A, \eta \rangle
\end{aligned}$$

Turunan variasi dari  $A(\psi(x))$  adalah  $\delta A(\psi(x)) = x^3 \cos(\psi(x)) + 5\psi(x)^4 - \partial_{xx} u(x)$

## 2. Light rays, Fermat's principle.

According to Fermat, the trajectory of a light ray between two points is such that the time required time is as small as possible. The propagation speed of light depends on material properties, which is expressed by  $c_0/n$  where  $c_0$  is the speed in vacuum (which is maximal), and  $n>1$  is the so called index of refraction, characteristic for the material.

For trajectories, for simplicity described as graphs of functions  $x \rightarrow y(x)$ . The total time between points is

$$\int_D n(x, y) \sqrt{1 + y_x^2} \, dx \quad \dots (*)$$

this is also often called the optical pathlength. Note that this functional can also be given very different interpretations, depending on the meaning of  $n$ .

- a. Write down the Euler-Lagrange equations.
- b. Determine the optimal trajectory in case  $n$  does not depend on  $x$  explicitly.

Then use 'energy conservation' to study trajectories.

- c. Consider the special cases  $n = y$  and  $n = \frac{1}{y}$  for which the trajectories can be expressed explicitly.

Penyelesaian:

a. Misalkan  $A(y) = \int_D n(x, y) \sqrt{1 + y_x^2} dx$

Turunan variasi dari  $A(y)$ :

$$\begin{aligned}
\frac{\partial}{\partial \varepsilon} A(y + \varepsilon \eta) \Big|_{\varepsilon=0} &= \frac{\partial}{\partial \varepsilon} \int_D n(x, y + \varepsilon \eta) \sqrt{1 + \left( \frac{\partial y + \varepsilon \eta}{\partial x} \right)^2} dx \Big|_{\varepsilon=0} \\
&= \int_D \frac{\partial n(x, y + \varepsilon \eta)}{\partial y} \eta \sqrt{1 + \left( \frac{\partial y + \varepsilon \eta}{\partial x} \right)^2} dx + \\
&\quad \int_D n(x, y + \varepsilon \eta) \frac{\frac{\partial y + \varepsilon \eta}{\partial x}}{\sqrt{1 + \left( \frac{\partial y + \varepsilon \eta}{\partial x} \right)^2}} dx \\
&= \int_D \left( \frac{\partial n(x, y)}{\partial y} \sqrt{1 + \left( \frac{\partial y}{\partial x} \right)^2} \right) \eta dx + \int_D n(x, y) \frac{\frac{\partial y}{\partial x}}{\sqrt{1 + \left( \frac{\partial y}{\partial x} \right)^2}} \eta dx \\
&= \int_D \left( \frac{\partial n(x, y)}{\partial y} \sqrt{1 + \left( \frac{\partial y}{\partial x} \right)^2} \right) \eta dx + \frac{n(x, y) \frac{\partial y}{\partial x}}{\sqrt{1 + \left( \frac{\partial y}{\partial x} \right)^2}} \eta \Big|_D - \int_D \frac{\partial}{\partial x} \left( \frac{n(x, y) \frac{\partial y}{\partial x}}{\sqrt{1 + \left( \frac{\partial y}{\partial x} \right)^2}} \right) \eta dx \\
&= \frac{n(x, y) \frac{\partial y}{\partial x}}{\sqrt{1 + \left( \frac{\partial y}{\partial x} \right)^2}} \eta \Big|_D + \int_D \left( \frac{\partial n(x, y)}{\partial y} \sqrt{1 + \left( \frac{\partial y}{\partial x} \right)^2} - \frac{\partial}{\partial x} \left( \frac{n(x, y) \frac{\partial y}{\partial x}}{\sqrt{1 + \left( \frac{\partial y}{\partial x} \right)^2}} \right) \right) \eta dx
\end{aligned}$$

Persamaan Euler-Lagrange :

$$\begin{aligned}
&\frac{\partial n(x, y)}{\partial y} \sqrt{1 + \left( \frac{\partial y}{\partial x} \right)^2} - \frac{\partial}{\partial x} \left( \frac{n(x, y) \frac{\partial y}{\partial x}}{\sqrt{1 + \left( \frac{\partial y}{\partial x} \right)^2}} \right) = 0 \\
\Leftrightarrow &\partial_y n \sqrt{1 + \left( \frac{\partial y}{\partial x} \right)^2} - \frac{\frac{\partial_x n}{\partial y} \frac{\partial y}{\partial x} + \frac{\partial_y n}{\partial x} \frac{\partial y}{\partial x}^2}{\sqrt{1 + \left( \frac{\partial y}{\partial x} \right)^2}} - n \frac{\partial_{xx} y}{\sqrt{1 + \left( \frac{\partial y}{\partial x} \right)^2}} + n \frac{\frac{\partial y}{\partial x}^2 \frac{\partial_{xx} y}{\partial x}}{\left( 1 + \frac{\partial y}{\partial x}^2 \right)^{3/2}} = 0
\end{aligned}$$

Persamaan terakhir disederhanakan menjadi:

$$\partial_y n \left( 1 + (\partial_x y)^2 \right) - n \partial_{xx} y - \partial_x n \partial_x y \left( 1 + (\partial_x y)^2 \right) = 0$$

b. Dalam kasus  $n(x, y)$  tidak begantung secara eksplisit pada  $x$ , maka  $\frac{\partial n(x, y)}{\partial x} = 0$ ,

sehingga persamaan Euler-Lagrange menjadi

$$\partial_y n \sqrt{1 + \frac{\partial_x y^2}{\partial_x y}} - \frac{\partial_y n \partial_x y^2}{\sqrt{1 + \frac{\partial_x y^2}{\partial_x y}}} - n \frac{\partial_{xx} y}{\sqrt{1 + \frac{\partial_x y^2}{\partial_x y}}} + n \frac{\partial_x y^2 \partial_{xx} y}{\left(1 + \frac{\partial_x y^2}{\partial_x y}\right)^2} = 0$$

$$\Leftrightarrow \partial_y n \left(1 + (\partial_x y)^2\right) - n \partial_{xx} y = 0$$

$$\Leftrightarrow \partial_y n + \partial_y n \partial_x y^2 - n \partial_{xx} y = 0$$

*Trajectory* (solusi) dari persamaan diferensial di atas, berupa kurva kuadratik.

c. – Untuk kasus  $n = y(x)$

Berdasarkan no 8.a maka persamaan Euler-Lagrange untuk kasus ini adalah:

$$\sqrt{1 + \frac{\partial_x y^2}{\partial_x y}} - \frac{\partial_x y^2}{\sqrt{1 + \frac{\partial_x y^2}{\partial_x y}}} - \frac{y \partial_{xx} y}{\sqrt{1 + \frac{\partial_x y^2}{\partial_x y}}} + \frac{y \partial_x y^2 \partial_{xx} y}{\left(1 + \frac{\partial_x y^2}{\partial_x y}\right)^2} = 0$$

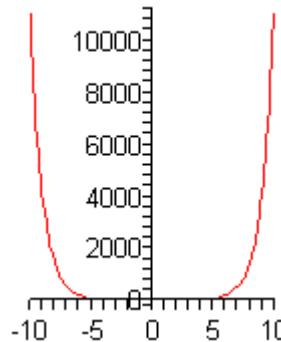
Dengan menggunakan program Maple (terdapat pada lampiran) solusi dari persamaan differensial tersebut adalah

$$y(x) = \frac{1 + e^{c_1 x^2} (e^{c_2 x})^2}{2c_1 e^{c_1 x^2} e^{c_2 x}}, \text{ dimana } c_1 \text{ dan } c_2 \text{ adalah konstanta}$$

Dengan mengambil nilai awal  $y(0) = 1$  dan  $y'(0) = 0$  diperoleh solusi :

$$y(x) = \frac{1 + e^{-x^2}}{2e^{-x}}$$

dan gambar solusi dari persamaan diferensial tersebut (*trajectory*) adalah



$\therefore$  Light Rays di dalam medium dengan fungsi  $n(x, y) = y(x)$  adalah berupa kurva kuadratik.

- Untuk kasus  $n = \frac{1}{y(x)}$

Berdasarkan no 8.a maka persamaan Euler-Lagrange untuk kasus ini :

$$\frac{1}{y^2} \sqrt{1 + \dot{y}^2} + \frac{\ddot{y}}{y^2 \sqrt{1 + \dot{y}^2}} - \frac{\ddot{y}_{xx} y}{y \sqrt{1 + \dot{y}^2}} + \frac{\dot{y}^2 \ddot{y}_{xx} y}{y^2 (1 + \dot{y}^2)^{3/2}} = 0$$

Dengan menggunakan program Maple (terdapat pada lampiran) solusi dari persamaan differensial tersebut adalah

$$y_1(x) = \sqrt{-x^2 - 2c_1 x + 2c_2}, \text{ dan } y_2(x) = -\sqrt{-x^2 - 2c_1 x + 2c_2},$$

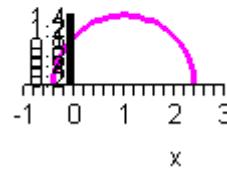
$c_1$  dan  $c_2$  adalah konstanta.

Dengan mengambil nilai awal  $y(0) = 0$  dan  $y_x(0) = 1$  diperoleh solusi:

$$y(x) = \sqrt{-x^2 + 2x + 1}$$

Gambar trajectory dari persamaan diferensial tersebut adalah

lintasan sinar



*Light rays* di dalam medium dengan  $n(x, y) = \frac{1}{y(x)}$  berupa kurva parabola (kuadratik).

### 3. Boussinesq type of equations.

Surface waves (in one horizontal direction  $x$ ) that decay at infinity ( $|x| \rightarrow \pm\infty$ ) can be described in terms of the wave height  $p(x, t)$  and a velocity  $u(x, t)$  in the following form (a Hamiltonian system):

$$\begin{aligned} \partial_t u &= \partial_x \nabla H(u, p), \\ \partial_t p &= -\partial_x \nabla_u H(u, p) \end{aligned} \quad \dots(i)$$

For a suitable functional (the Hamiltonian)  $H(u, p)$ .

- Describe the equations in full detail when the Hamiltonian is given by the following functional

$$H(u, p) = \int \left\{ \frac{1}{2} gp^2 + \frac{1}{2} (h + p) \left( u^2 - \frac{1}{3} u_x^2 \right) \right\} dx$$

(This set of equations are the linearized equations).

- b. In another case (shallow water, no dispersion, but nonlinear), the equations are of the form

$$\begin{aligned}\partial_t u &= \partial_x \left\{ gp + \frac{1}{2} u^2 \right\}, \\ \partial_t p &= -\partial_x (u + \beta pu)\end{aligned}$$

Where  $\beta$  is constant. Determine the value of  $\beta$  such that this system of equations is a Hamiltonian system of the form (3,4) given above.

- c. Show that the equations have the horizontal momentum as constant of the motion

$$\int u(x)p(x)dx$$

Penyelesaian:

a. Hamiltonian:  $H(u, p) = \int \left\{ \frac{1}{2} gp^2 + \frac{1}{2} (h + p) \left( u^2 - \frac{1}{3} u_x^2 \right) \right\} dx$

- Akan ditentukan  $\delta_u H(u, p)$

$$\begin{aligned}\frac{\partial}{\partial \varepsilon} H(u + \varepsilon \eta, p) \Big|_{\varepsilon=0} &= \frac{\partial}{\partial \varepsilon} \int \frac{1}{2} gp^2 + \frac{1}{2} (u + p \left[ u + \varepsilon \eta - \frac{1}{3} \frac{\partial u + \varepsilon \eta}{\partial x} \right]) dx \Big|_{\varepsilon=0} \\ &= \int (u + p \eta) dx - \int \frac{1}{3} (u + p \eta_x u_x p \eta_x) dx \\ &= \int (u + p \eta) dx - \frac{1}{3} (u + p \eta_x u \eta_x) \Big|_{x=a}^b + \int \frac{1}{3} [h_x h + \partial_x p \partial_x u + (u + p \partial_{xx} u) \eta] dx \\ &= -\frac{1}{3} (u + p \eta_x u \eta_x) \Big|_{x=a}^b + \int \left[ (u + p \eta) + \frac{1}{3} (h_x h + \partial_x p \partial_x u + \frac{1}{3} (u + p \partial_{xx} u) \eta) \right] dx\end{aligned}$$

maka  $\delta_u H(u, p) = (u + p \eta) + \frac{1}{3} (h_x h + \partial_x p \partial_x u + \frac{1}{3} (u + p \partial_{xx} u) \eta)$

- Akan ditentukan  $\delta_p H(u, p)$

$$\begin{aligned}\frac{\partial}{\partial \varepsilon} H(p + \varepsilon \eta) \Big|_{\varepsilon=0} &= \frac{\partial}{\partial \varepsilon} \int \frac{1}{2} g \Phi + \varepsilon \eta^2 + \frac{1}{2} \Phi + p + \varepsilon \eta \left( u^2 - \frac{1}{3} \Phi_x u^2 \right) dx \Big|_{\varepsilon=0} \\ &= \int gp \eta dx - \int \frac{1}{2} \eta \left( u^2 - \frac{1}{3} \Phi_x u^2 \right) dx \\ &= \int \left[ gp + \frac{1}{2} \left( u^2 - \frac{1}{3} \Phi_x u^2 \right) \right] \eta dx\end{aligned}$$

$$\text{Maka } \delta_p H(u, p) = gp + \frac{1}{2} \left( u^2 - \frac{1}{3} \Phi_x u^2 \right)$$

Sehingga sistem (i) menjadi

$$\begin{aligned}\partial_t u &= \partial_x \left\{ gp + \frac{1}{2} \left( u^2 - \frac{1}{3} \Phi_x u^2 \right) \right\}, \\ \partial_t p &= -\partial_x \left\{ \Phi + p \cancel{u} + \frac{1}{3} \Phi_x h + \partial_x p \cancel{\partial}_x u + \frac{1}{3} \Phi + p \cancel{\partial}_{xx} u \right\}\end{aligned}$$

Yang ekuivalen dengan :

$$\begin{aligned}\partial_t u &= g \partial_x p + u \partial_x u - \frac{1}{3} \partial_x u - \partial_{xx} u, \\ \partial_t p &= -\partial_x u \cancel{h} + p \cancel{-} \partial_x h + \partial_x p \cancel{u} - \frac{1}{3} \partial_{xx} h + \partial_{xx} p \cancel{\partial}_x u - \frac{2}{3} \partial_x h + \partial_x p \cancel{\partial}_{xx} u \\ &\quad - \frac{1}{3} \cancel{h} + p \cancel{\partial}_{xxx} u\end{aligned}$$

- b. Akan ditentukan  $\beta$  sehingga  $\delta_p H(u, p) = gp + \frac{1}{2} u^2$  dan  $\delta_u H(u, p) = u + \beta p u$

$$\begin{aligned}\langle \delta_p H(u, p), \eta \rangle &= \int \left( gp + \frac{1}{2} u^2 \right) \eta dx \\ &= \int gp \eta + \frac{1}{2} u^2 \eta dx\end{aligned}$$

Untuk memperoleh  $H(u, p)$ , integralkan  $\int gp + \frac{1}{2} u^2 dx$  terhadap  $p$ .

$$\begin{aligned}H(u, p) &= \int \int \left( gp + \frac{1}{2} u^2 \right) dx dp \\ &= \int \frac{1}{2} gp^2 + \frac{1}{2} u^2 p + j(u) dx, \quad \text{dimana } j(u) \text{ adalah fungsi dari } u\end{aligned}$$

$$\begin{aligned}
\frac{\partial}{\partial \varepsilon} H(u + \varepsilon\eta, p) \Big|_{\varepsilon=0} &= \frac{\partial}{\partial \varepsilon} \int \frac{1}{2} gp^2 + \frac{1}{2} (u + \varepsilon\eta)^2 p - j(u + \varepsilon\eta) dx \Big|_{\varepsilon=0} \\
&= \int u\eta p - \partial_u j \eta dx \\
&= \int [p - \partial_u j] \eta dx
\end{aligned}$$

Maka  $\delta_u H(u, p) = up + \partial_u j$

Karena diketahui  $\delta_u H(u, p) = \beta pu + u$  maka  $\beta = 1$  dan  $j(u) = \frac{1}{2}u^2$

$\therefore$  Diperoleh  $\beta = 1$  dengan fungsi Hamiltonian  $H(u, p) = \int \frac{1}{2} gp^2 + \frac{1}{2} u^2 (p+1) dx$

### c. Horizontal momentum

Berdasarkan no.8b diperoleh bahwa

$$\partial_t u = gp_x + uu_x$$

$$\partial_t p = -u_x - up_x - pu_x$$

Akan ditunjukkan  $\frac{d}{dt} \int u(x, t) p(x, t) dx = 0$

$$\begin{aligned}
\frac{d}{dt} \int u(x, t) p(x, t) dx &= \int \frac{d(u(x, t)p(x, t))}{dt} dx \\
&= \int [\partial_t u] p + u [\partial_t p] dx \\
&= \int [gp_x + uu_x] p + [u_x - up_x - pu_x] dx \\
&= \int gp_x p - uu_x p - p_x u^2 dx \\
&= \frac{1}{2} gp^2 - \frac{1}{2} u^2 - \int p_x u^2 dx
\end{aligned}$$