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# KINEMATIKA GELOMBANG

## TOPIK 2

## KULIAH GELOMBANG OPTIK

ANDHY SETIAWAN

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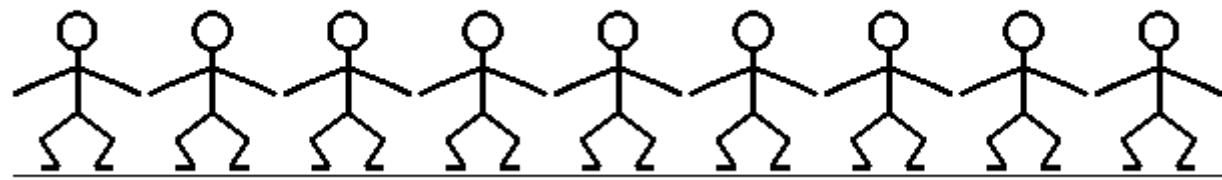


# SUB POKOK BAHASAN

- PERSAMAAN DIFFERENSIAL GELOMBANG
- SOLUSI PERSAMAAN GELOMBANG
- SUPERPOSISI DUA GELOMBANG

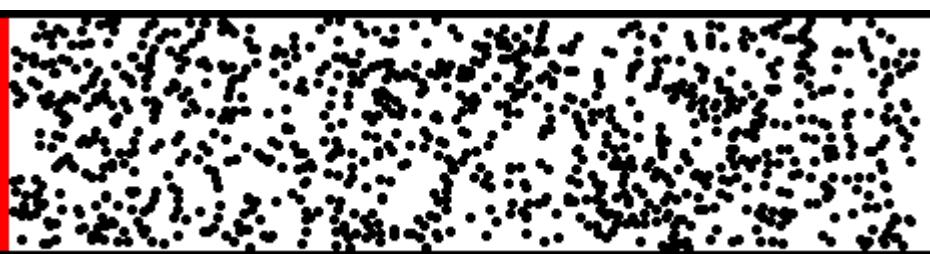
PENGANTAR

# ILUSTRASI PERAMBATAN PULSA



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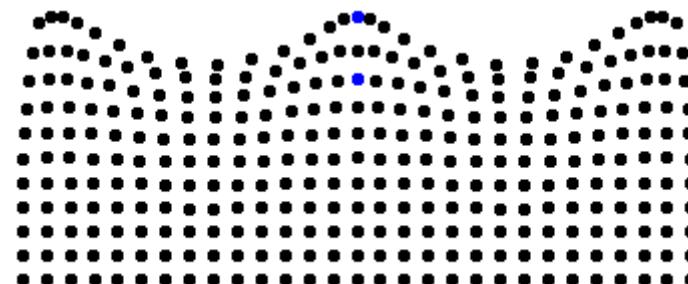
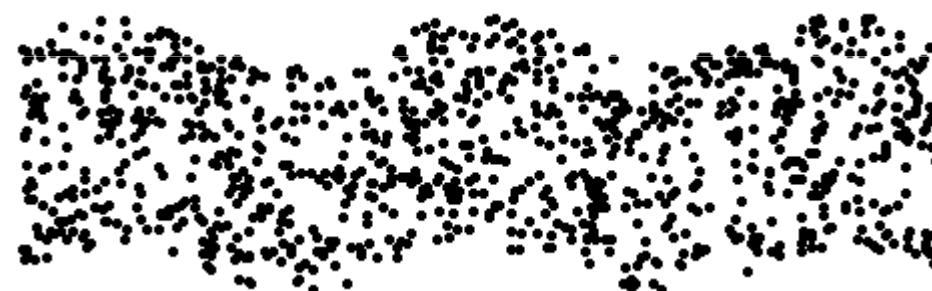
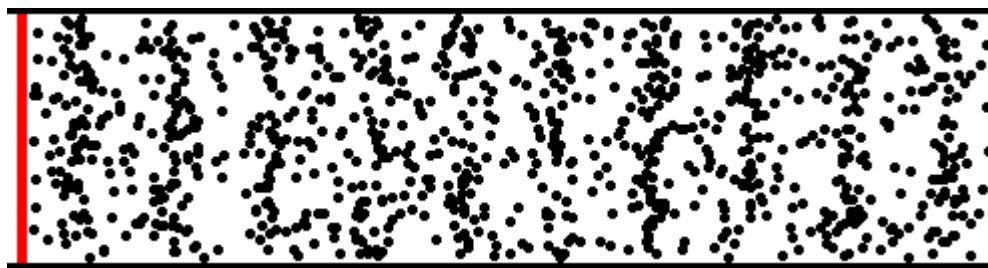
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# PENGANTAR ILUSTRASI PERAMBATAN GELOMBANG



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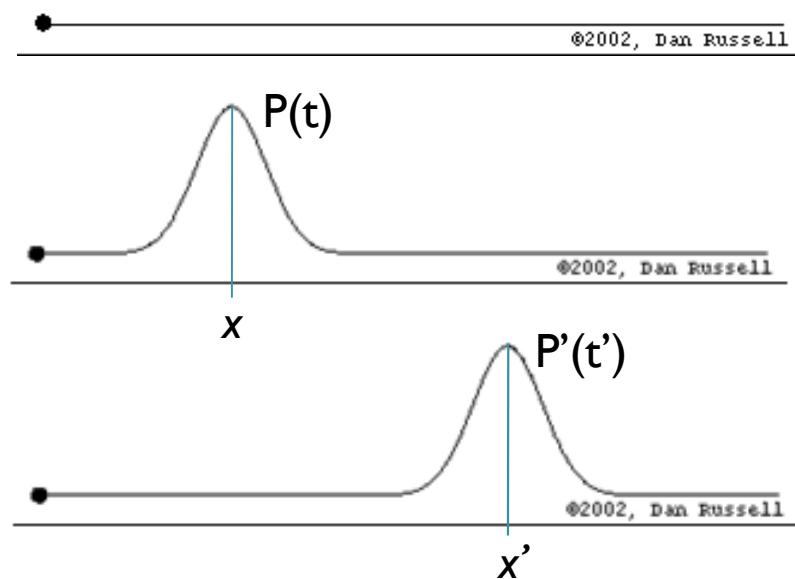


## PERSAMAAN DIFFERENSIAL GELOMBANG arah rambat dan sudut fase

- Sistem osilasi  $\psi(t)$
  - fungsi gelombang  $\psi(x, t)$  atau  $\psi(r, t)$
  - Tinjau: merambat arah  $x$ , kecepatan konstan  $v$ .  $\rightarrow \psi(x, t) = f(x \pm vt)$   
 $\psi(x, t) = f(\phi)$ , dengan  $\phi = x \pm vt$
- $\phi = \text{sudut fase}$

# PERSAMAAN DIFFERENSIAL GELOMBANG

## arah rambat dan sudut fase



Sudut fase titik P :  $\phi = x - vt$

Setelah t' :  $\phi' = x' - vt'$

$$\begin{aligned}\phi &= \phi', \\ x - vt &= x' - vt'\end{aligned}$$

$$x - vt = x + \Delta x - v(t + \Delta t)$$

$$0 = \Delta x - v \Delta t$$

$$\Delta x = v \Delta t$$

Maka  $\Delta x > 0$ , sehingga :

sudut fase  $\phi = x - vt \rightarrow$  arah rambat ke kanan

sudut fase  $\phi = x + vt \rightarrow$  arah rambat ke kiri (**coba buktikan**)

# PERSAMAAN DIFFERENSIAL GELOMBANG

## penurunan persamaan

- $\phi = x \pm vt$  konstan  $\rightarrow$  kedudukan setiap titik yang sama

Kecepatan fase

$$\frac{d\phi}{dt} = 0 \rightarrow \frac{d(x \pm vt)}{dt} = 0 \rightarrow \frac{dx}{dt} \pm v = 0 \rightarrow v = \mp \frac{dx}{dt}$$

- Perubahan fungsi terhadap  $x$  dan  $t$

$$\frac{\partial \psi}{\partial x} = \frac{\partial \psi}{\partial \phi} \frac{\partial \phi}{\partial x} = \frac{\partial \psi}{\partial \phi} \rightarrow \frac{\partial \psi}{\partial \phi} = \frac{\partial \psi}{\partial x} \rightarrow \frac{\partial \psi}{\partial x} \mp \frac{1}{v} \frac{\partial \psi}{\partial t} = 0$$

$$\frac{\partial \psi}{\partial t} = \frac{\partial \psi}{\partial \phi} \frac{\partial \phi}{\partial t} = \pm v \frac{\partial \psi}{\partial \phi} \rightarrow \frac{\partial \psi}{\partial \phi} = \pm \frac{1}{v} \frac{\partial \psi}{\partial t}$$

# PERSAMAAN DIFFERENSIAL GELOMBANG

## penurunan persamaan

- Turunan kedua terhadap  $x$  dan  $t$

$$\frac{\partial^2 \psi}{\partial x^2} = \frac{\partial}{\partial x} \left( \frac{\partial \psi}{\partial x} \right) = \frac{\partial}{\partial x} \left( \frac{\partial \psi}{\partial \phi} \right) = \frac{\partial}{\partial \phi} \left( \frac{\partial \psi}{\partial x} \right) = \frac{\partial}{\partial \phi} \left( \frac{\partial \psi}{\partial \phi} \right) = \frac{\partial^2 \psi}{\partial \phi^2}$$

$$\frac{\partial^2 \psi}{\partial t^2} = \frac{\partial}{\partial t} \left( \frac{\partial \psi}{\partial t} \right) = \frac{\partial}{\partial t} \left( \pm v \frac{\partial \psi}{\partial \phi} \right) = (\pm v) \frac{\partial}{\partial \phi} \left( \frac{\partial \psi}{\partial t} \right) = (\pm v) \frac{\partial}{\partial \phi} \left( \pm v \frac{\partial \psi}{\partial \phi} \right) = v^2 \frac{\partial^2 \psi}{\partial \phi^2}$$

$$\frac{\partial^2 \psi}{\partial \phi^2} = \frac{\partial^2 \psi}{\partial x^2} \quad \boxed{\frac{\partial^2 \psi}{\partial x^2} - \frac{1}{v^2} \frac{\partial^2 \psi}{\partial t^2} = 0} \quad \frac{\partial^2 \psi}{\partial \phi^2} = \frac{1}{v^2} \frac{\partial^2 \psi}{\partial t^2}$$

$$\nabla^2 \psi - \frac{1}{v^2} \frac{\partial^2 \psi}{\partial t^2} = 0$$

Merupakan ungkapan gelombang datar  
(Front wave berupa bidang datar)

Untuk koordinat bola

$$\nabla^2 \psi = \frac{1}{r^2} \frac{\partial}{\partial r} r^2 \frac{\partial \psi}{\partial r} \quad \rightarrow \quad \frac{\partial^2 \psi}{\partial r^2} + \frac{2}{r} \frac{\partial \psi}{\partial r} - \frac{1}{v^2} \frac{\partial^2 \psi}{\partial t^2} = 0 \quad (\text{Buktikan})$$

# PERSAMAAN DIFFERENSIAL GELOMBANG

## prinsip superpoisi

- Jika  $\psi_1$  dan  $\psi_2$  solusi dari pers. Gelombang, maka berlaku:

$$\frac{\partial^2 \psi_1}{\partial x^2} - \frac{1}{v^2} \frac{\partial^2 \psi_1}{\partial t^2} = 0$$

$$\frac{\partial^2 \psi_2}{\partial x^2} - \frac{1}{v^2} \frac{\partial^2 \psi_2}{\partial t^2} = 0$$

dijumlahkan

$$\frac{\partial^2 \psi}{\partial x^2} - \frac{1}{v^2} \frac{\partial^2 \psi}{\partial t^2} = 0$$

$$\frac{\partial^2(\psi_1 + \psi_2)}{\partial x^2} - \frac{1}{v^2} \frac{\partial^2(\psi_1 + \psi_2)}{\partial t^2} = 0$$

Jadi  $(\psi_1 + \psi_2)$  merupakan solusi dari pers. Gelombang juga

Prinsip superposisi

# SOLUSI PERSAMAAN GELOMBANG

Solusi paling sederhana dari persamaan :  $\frac{\partial^2 \Psi}{\partial x^2} - \frac{1}{v^2} \frac{\partial^2 \Psi}{\partial t^2} = 0$  adalah

$$\Psi(x,t) = \Psi_0 \cos k(x-vt) \longrightarrow \Psi_0 = \Psi_{\text{maks}}$$

$k$  = bilangan gelombang/vektor gelombang (menunjukan arah rambat gelombang)

$$\Psi(x,t) = \Psi_0 \cos k(x-vt)$$

$k$  = frekuensi spatial

$\omega$  = frekuensi temporal

$$\Psi(x,t) = \Psi_0 \cos (kx - kvt)$$

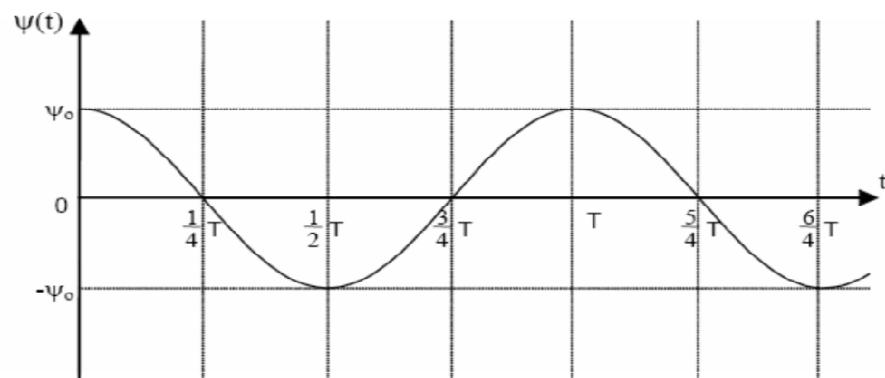
$$\Psi(x,t) = \Psi_0 \cos (kx - \omega t)$$

$$k = \frac{2\pi}{\lambda}$$

$$\omega = \frac{2\pi}{T}$$

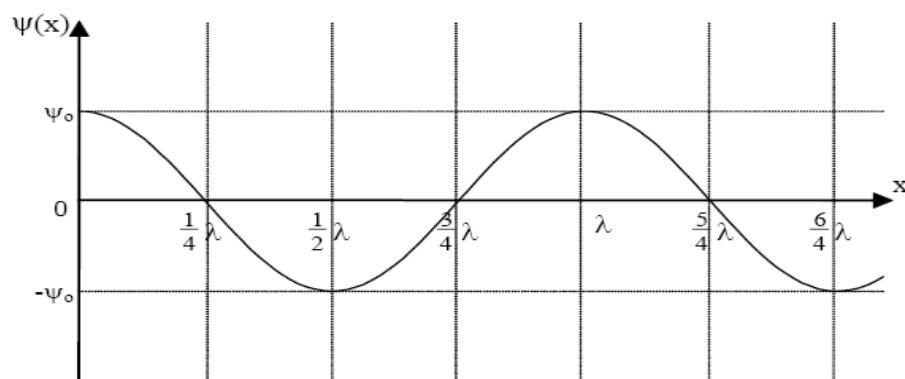
$\lambda$  = perioda spatial

$T$  = perioda temporal



Mengungkapkan  
pola eksitasi  
gelombang

Gelombang dalam sisi temporal



Mengungkapkan  
perambatan  
gelombang

Gelombang dalam sisi spatial

Sehingga solusi persamaan gelombang dapat pula diungkapkan dengan:

$$\Psi(x, t) = A \cos\left[2\pi\left(\frac{x}{\lambda} - \frac{t}{T}\right)\right]$$

# SUPERPOSISI DUA GELOMBANG

Misalkan dua buah gelombang dengan arah getar pada bidang yang sama, masing-masing frekuensinya  $\omega_1$  dan  $\omega_2$  serta bilangan gelombangnya  $k_1$  dan  $k_2$

$$\psi_1(x,t) = A \cos(k_1 x - \omega_1 t) \quad \text{dan} \quad \psi_2(x,t) = A \cos(k_2 x - \omega_2 t)$$

Hasil superposisinya adalah:

$$\Psi(x,t) = A[\cos(k_1 x - \omega_1 t) + \cos(k_2 x - \omega_2 t)]$$

$$\begin{aligned}\Psi(x,t) \\ = 2A \left[ \cos\left\{\frac{(k_1 - k_2)x - (\omega_1 - \omega_2)t}{2}\right\} \cos\left\{\frac{(k_1 + k_2)x - (\omega_1 + \omega_2)t}{2}\right\} \right]\end{aligned}$$

$$\Delta k = k_1 - k_2 \quad \Delta\omega = \omega_1 - \omega_2 \quad \text{Maka:}$$

$$\Psi(x, t)$$

$$= 2A \left[ \cos \left\{ \frac{\Delta kx - \Delta \omega t}{2} \right\} \cos \left\{ \frac{(2k_1 - \Delta k)x - (2\omega_1 - \Delta \omega)t}{2} \right\} \right]$$

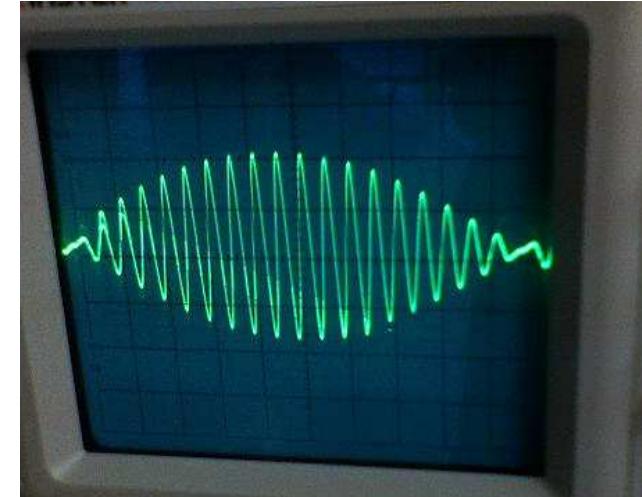
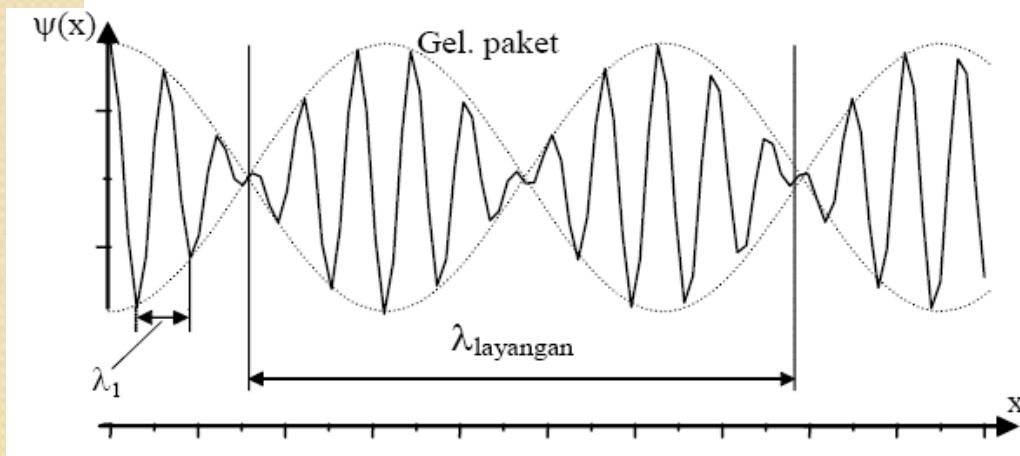
Untuk sisi spasial (dengan t konstan, misal t=0)

$$\Psi(x, t) = 2A \left[ \cos \frac{\Delta kx}{2} \cos \frac{(2k_1 - \Delta k)x}{2} \right]$$

$\Delta k$  sangat kecil, sehingga  $2k_1 - \Delta k \approx 2k_1$

$$\Psi(x, t) = 2A \left[ \cos \frac{\Delta kx}{2} \cos k_1 x \right]$$

Bila kita gambarkan hasil superposisinya, maka :

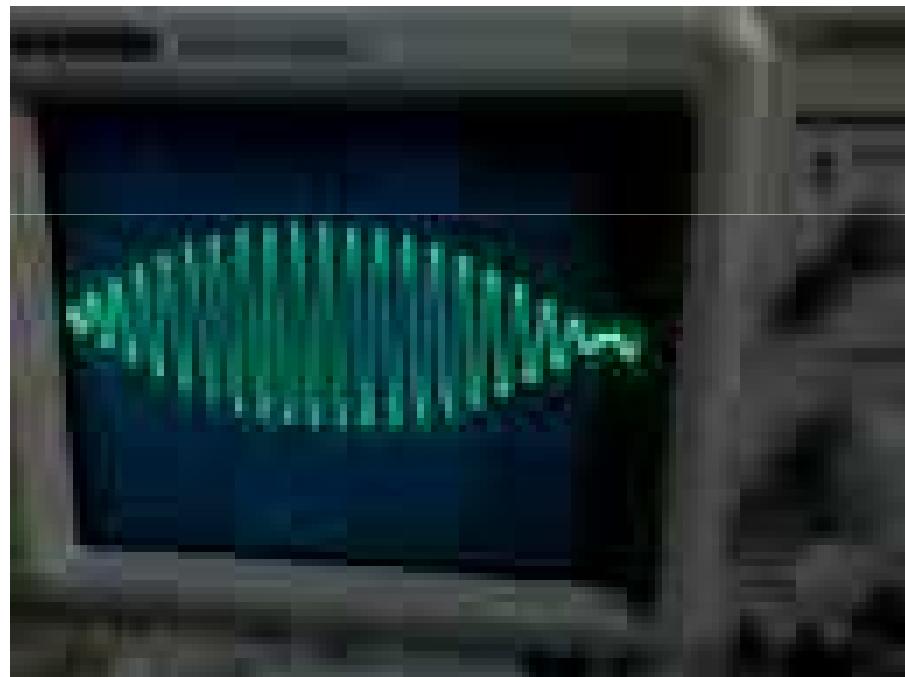
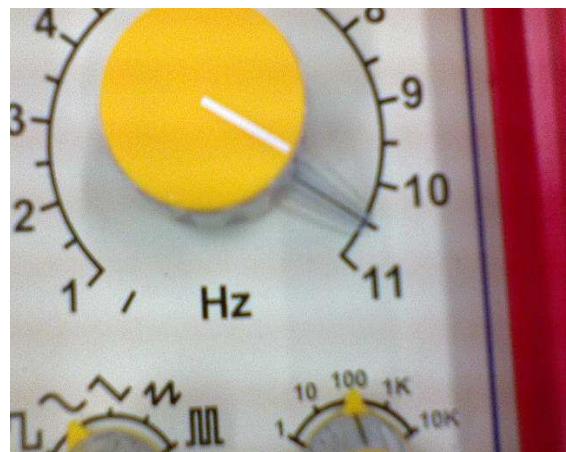


Hasil superposisi kedua gelombang dengan perbedaan frekuensi yang kecil ini disebut *layangan*, hasilnya berupa gelombang paket yang terselubung (envelope), dan kecepatan gelombang paket ini disebut dengan kecepatan group.

$$\text{Kecepatan fase: } v = \frac{\omega_1}{k_1}$$

$$\text{Kecepatan group: } v_g = \frac{\Delta\omega}{\Delta k} = \frac{d\omega}{dk} = \frac{d(kv)}{dk} = v + k \frac{dv}{dk}$$

# Layangan



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# SUPERPOSI DUA GELOMBANG arah getar saling tegak lurus

Tinjauan dua gelombang dengan frekuensi yang sama dan arah getar yang tegak lurus: Misal arah getarnya Y dan Z:

$$\psi_y(t) = A_1 \sin(\omega t + \varphi_1)$$

$$\psi_z(t) = A_2 \sin(\omega t + \varphi_2)$$

$$\frac{\Psi_y}{A_1} \cos \varphi_2 - \frac{\Psi_z}{A_2} \cos \varphi_1$$

$$= \cos(\omega t) \{ \cos(\varphi_2) \sin(\varphi_1) - \cos(\varphi_1) \sin(\varphi_2) \}$$

Superposisi keduanya menghasilkan:

$$\begin{aligned} & -\frac{\Psi_y}{A_1} \sin \varphi_2 + \frac{\Psi_z}{A_2} \sin \varphi_1 \\ & = \sin(\omega t) \{ \cos(\varphi_2) \sin(\varphi_1) - \cos(\varphi_1) \sin(\varphi_2) \} \end{aligned}$$

Kuadratkan kedua persamaan, kemudian dijumlahkan, menghasilkan:

$$\sin^2(\delta) = \left(\frac{\Psi_y}{A_1}\right)^2 + \left(\frac{\Psi_z}{A_2}\right)^2 - \frac{2\Psi_y\Psi_z}{A_1A_2} \cos(\delta)$$

Dengan beda sudut fase:  $\delta = \varphi_1 - \varphi_2$

Persamaan ini merupakan persamaan umum elips, karena itu superposisinya disebut terpolarisasi elips.

Untuk beberapa kasus khusus, yaitu:  $\delta = \pi/2, 3\pi/2, 5\pi/2.....$ , persamaanya jadi:

$$\left(\frac{\Psi_y}{A_1}\right)^2 + \left(\frac{\Psi_z}{A_2}\right)^2 = 1$$

Terjadi polarisasi elips putar kanan, dan bila amplitudo kedua gelombang sama ( $A_1=A_2$ ), maka superposisinya terpolarisasi lingkaran putar kanan.

Bila:  $\delta = 0, 2\pi, 4\pi, ....$  Persamaan menjadi:  $\Psi_y = \frac{A_1}{A_2} \Psi_z$

Terjadi polarisasi linear